$: 10013806(2002)05035004$ 

# Pulsed laser performance of Nd GdCOB and Cr: Nd GdCOB crystals

H ou Xueyuan, H uang Jungang , Sun Yuming , Li Yuf ei , Liu Enquan ( Department of Optics, Shandong University, Jinan, 250100)

Abstract:  $Ca_4GdO(BO_3)_3(GdCOB)$  is a new self frequency doubling (SFD) crystal. Using x enon lamp as pump sources, we have realized the free run from 1061nm to 530. 5nm in SFD Cr- Nd- GdCOB( bi doped) crystal and Nd-GdCOB(unidoped) crystal. The threshold energies for unidoped crystal and bidoped crystal are 1.0J and 0.92J, respectively; and the maximum output energies of green laser for two kinds of crystals are 1. 96mJ and 2. 46mJ, respectively. By using pulsed dye laser as pump source, we have obtained SFD red laser at 655nm and its fundamental frequency laser at 1331nm as well as achieved SFD green laser at 530. 5nm from a Nd-GdC0B crystal.

Key words: self frequency doubling; phase matching; unidoped; dye laser; x enon lamp

#### Nd: GdCOB Cr: 1 Nd: GdCOB

侯学元 黄俊刚 孙渝明 李宇飞 刘恩泉  $($ , , 250100)

:  $Ca_4GdO(BO_3)$  3(  $GdCOB$ ) Cr-Nd- $1061$ nm~ 530. 5nm 1. OJ 0.  $92J,$  1.  $1.96mJ$  $2.46mJ$ Nd GdCOB GdCOB 1331nm 655nm , 530. 5nm : ; ; ; ; : O734 : A

### Introduction

Bluegreen region lasers have potential applications in medicine, high brightness display, undersea applications, high density data storage and so on. Self-frequency-doubling (SFD) crystal can be used to realize the compactness and miniaturization of this kind of laser. One of the famous SFD crystals, NYAB, has large nonlinear coefficient, but its difficult growth with large size and homogeneity problem limits its further application and development<sup>[1,2]</sup>.

Nd: Ca4GdO(BO3)3(Nd GdCOB) is a new SFD crystal. It belongs to monoclinic biaxial crystal system. Its cell parameters are  $a = 0.8095(7)$  nm,  $b= 1.60018(6)$  nm,  $c= 0.3558(8)$  nm,  $\beta= 101.26^{\circ}$ . The transmission spectrum ( $350 \text{nm} \sim 2700 \text{nm}$ ) of Nd-GdCOB crystal is little narrow er than that of

 $: 2001 - 08 + 17;$   $: 2001 - 09 + 26$ 

BBO crystal ( $200 \text{nm} \sim 2600 \text{nm}$ ); its effective nonlinear coefficient is about 1. 0pm/ V, lower than  $BBO (2pm/ V)$ , higher than  $LBO (0. 82pm/ V)$  and equivalent to NYAB. Nd-GdCOB crystal has the advantages of high damage threshold (up to  $1$ GW/ cm 2 ) and easy growth w ith large size. GdCOB crystal can realize type I phase matching in the spectrum ranges  $720nm \sim 1500nm$  and  $840nm \sim 2000nm$ , including red, green, blue frequency doubling band. Its three primary fluorescence spectrums are: 936nm, 1061nm, 1331nm, and they can be frequency-doubled to 468nm ( blue) , 530nm ( green) , 665. 5nm ( red) , respectively  $[3-6]$ .

Nd: GdCOB

Using xenon lamp as pump source, w e realize the free run from 1061nm to 530. 5nm in SFD Nd: GdCOB crystal and Cr: Nd: GdCOB crystal, respectively . T he experimental result show s that the output energy of Cr: Nd: GdCOB crystal is 25% higher than that of Nd: GdCOB crystal. Moreover, using  $Q$ -switched pulsed dye laser to pump

<sup>:</sup> , , 1948

Nd-GdCOB crystal, we obtained not only SFD green laser output at 530. 5nm and its fundamental frequency laser output at 1061nm but also observed SFD red laser output at 665. 5nm and its fundamental frequency laser output at 1331nm. This paper reports the pulsed laser characteristics of Nd-GdCOB crystal and Cr-Nd-GdCOB crystal.

# 1 Laser performance of Nd: GdCOB crystal and Cr: Nd: GdCOB crystal under xenon lamp pumping

### 1. 1 Manufacture of crystal and transmission spectrum

 $\mathcal{O}30$ mm × 40mm Nd: GdCOB(7% Nd) and Cr: Nd: GdCOB (1% Cr, 7% Nd) crystals have been developed by high-frequency heating Czochralski method. The crystals have excellent optical quality with no bubble. By using computer analogue, we obtain the phase matching curve of GdCOB and then know that GdCOB crystal can only realize type  $\bar{I}$ phase matching at 1061nm and the optimal phase matching angles are  $\theta$ = 66.  $\mathcal{E}, \phi$ = 47. 4°<sup>[7,8]</sup>. Both



Fig. 1 The transmission spectra of Cr. Nd GdCOB and Nd GdCOB a– Cr: Nd GdCOB b– Nd GdCOB

Nd: GdCOB and Cr: Nd: GdCOB crystal samples are cut at type I phase matching and have a dimension of  $3mm \times 3mm \times 7mm$ . They are not AR-coated at 1061nm and 530. 5nm. T he transmission spectrums of two crystals measured with HITACHI U-3500 spectrometer at room temperature are shown in Fig. 1. Due to the sensitization of  $Cr^{3+}$ , when the w aveleng th is shorter than 500nm, the absorption of Cr. Nd. GdCOB crystal is apparently greater than that of Nd-GdCOB crystal. Betw een 300nm and 1000nm, the ratio of the normalized absorption areas of two crystals is 1.5:1, which shows that Cr: Nd: GdCOB crystal can absorb more xenon lamp pumped light and is more appropriately pumped by a xenon lamp.

Meanwhile, Fig. 1 shows that the doping of  $Cr^{3+}$  ion haven't increased the absorption of Cr: Nd: GdCOB crystal for SFD light at 530. 5nm.

#### 1. 2 Experiment

In this experiment, the leng th of resonator is 130mm and the dimension of xenon lamp is  $\mathfrak{G}_{mm} \times$ 40mm. The pump light is focused by a 40mm long, single silver-coated elliptical cylinder reflector. Both the plane concave cavity and plane plane cavity are used in the experiment. The output mirror  $M_2$  has high reflection (99.9%) at 1061nm and high transmission  $(92\%)$  at 530. 5nm. For plane concave cavity, the concave mirror  $M_1$ , which has a curvature radius of 250mm, has high reflection at 1061nm ( 99. 8% ) and 530. 5nm( 99. 2%) , and the crystal is placed near  $M_2$ . For plane plane cavity, the plane mirror  $M_1$  has high reflection at 1061nm (99. 8%) and 530.  $5\text{nm}(99.4\%)$ , and the crystal is placed in the middle of cavity.

The experimental results show that the conversion efficiency of SFD laser in the plane plane cavity is only  $50\%$  of that in the plane concave cavity. This is because plane concave cavity can increase the energy density of the oscillating fundamental frequency laser. In the experiment, the scientedr 362 energy meter is used to measure the energy of output laser, and a split prism is used to separate the fundamental frequency laser from SDF laser.

#### 1. 3 Experimental results and discussion

Fig. 2 gives the v ariation of output green laser energy versus pump energy for Cr: Nd: GdCOB



Fig. 2 Comparative green laser output performances of Cr. Nd GdCOB and Nd-GdCOB crystal under xenon lam p pumping

crystal and Nd. GdCOB crystal. For Cr. Nd. GdCOB crystal, w hen the pump energy is 10J, the output energy of g reen laser is 2. 46mJ, which corresponds to

an increase of 25% in comparison to that of Nd-GdCOB crystal. The figure shows that the output energies of tw o crystals have less difference w hen the pump energy is low er. But w ith increasing pump energy, the difference of the output energies of two crystals increases gradually. The reason may be that the pulsed xenon lamp mainly emits the continuous spectrum under a high electric current density. With increasing the electric current density, the increase rate of short wavelength is faster than that of longw avelength. The  $\mathrm{Cr}^{3+}$  of  $\mathrm{Cr}$ : Nd: GdCOB crystal mainly absorbs the light of shortw aveleng th, and quickly transfers the absorbed energy to the excited Nd 3+ ion. For the SFD green laser at 530. 5nm, the pump threshold energies of Cr. Nd. GdCOB and Nd. GdCOB are 0. 92J and 1. 00J, respectively.

### 2 Laser characteristics of Nd: GdCOB crystal pumped by dye laser

From Fig. 1b, we know that there are two strong absorption peaks at 595nm and 811nm, which are suitable for dye laser and LD pump, respectively.

#### 2. 1 Experiment

In this experiment, a Nd: GdCOB rod (7%  $\mathrm{Nd}^{3+}$  ,  $3\mathrm{mm} \times 3\mathrm{mm} \times 10\mathrm{mm}$  cut at type I phase matching is placed in the middle of a 30mm-long plane plane cavity. Neither of two faces of the crystal has been AR-coated at 1061nm and 530. 5nm. The mirror  $M_1$  has high reflection at 1061nm (99.8%) and at 530. 5nm (99.4%), and the output mirror  $M_2$ has high reflection (99.8%) at 1061nm and high transmission ( 92% ) at 530. 5nm. The pump light from Datachrom-5000  $Q$ -switch pulsed dye laser ( 595. 0nm, pulse w idth 10ns and repetition rate 10Hz) is focused by a cylindrical lens ( focal length 50mm) to the Nd-GdCOB crystal. T he dimension of light spot at the crystal is about  $1 \text{mm} \times 7 \text{mm}$ . The pump threshold energy is 1. 2mJ, and the pulse width of the doubling frequency output laser is 9ns. A split prism is required to separate the fundamental frequency lig ht w hen the doubling frequency output is measured.

#### 2. 2 Results and discussion

Fig . 3 shows the variation of output SFD laser

energy versus pumping energy w hen the pumping w avelength  $\lambda_p$  is 595. 0nm. When the pumping energy is 17. 5mJ, the SFD output energy is 1. 35mJ, corresponding to a conversion efficiency of 7. 7%. Because the faces of crystal are not AR-coated and the output energy of our dye laser is low , the max imum SFD output energy is 1. 35mJ in the experiment. We can expect that if the faces of crystal are coated AR film and the pump energy is increased, a higher conversion efficiency could be obtained. T he threshold energy is about 1. 2mJ and is low er in comparison w ith NYAB ( 2mJ) . This is because Nd: GdCOB crystal has good optical quality and its selfabsorption at 530. 5nm is only 50% of NYAB crystal's.



Fig. 3 T he output SFD laser energy versus pumping energy  $(\lambda_{p} = 595nm)$ 

The variety of the output 530. 5nm laser energ y versus pumping wavelength  $\lambda_p$  (591.5nm ~ 601. 0nm) is plotted in Fig. 4- From Fig . 4, w e can see that there is a narrow peak at 595. 0nm and a few small peaks w hich are well in agreement w ith the absorption spectrum.



Fig. 4 Variation of the output energy with the pump w avelength

## 3 1331nm laser performances of Nd-GdCOB crystal pumped by dye laser

#### 3. 1 Experiment

GdCOB is a monoclinic biaxial crystal with

significant polarized absorption property. According to the polarized transmission spectrum of Nd: GdCOB, we know that the absorption efficiencies are obv iously different at three polarization directions. We have processed two crystal samples with the same dimension of  $3 \text{mm} \times 3 \text{mm} \times 8 \text{mm}$ . The oscillating laser travels through one sample along  $x$ -ax is ( called  $x$ -ax is sample) and the other along  $y$ -axis ( called  $y$ -axis sample). The angles between the crystallophysic axes and the corresponding crystallographic axes are:  $(x, c) = 12^{\circ}, (z, a) = 23^{\circ},$  $(y, b) = -180^\circ$ . Because of the linear polarization of the pump dye laser, we can obtain three pump modes  $(E \parallel x, E \parallel \gamma, E \parallel z)$  by changing the angle between the polarization direction of the pump laser and the crystallophysic ax es of the crystal.

In the ex periment, Nd-GdCOB crystal is placed in the middle of a 40mm long plane plane cavity which is composed of reflective mirrors  $M_1$  and  $M_2$ . The  $M_1$  has a high reflectivity of 99. 8% at 1331. 0nm, and the transmittance of the output mirror  $M_2$  is variable. The 1331.0nm output energy is the hig hest w hen the w avelength of dye laser is 595. 5nm.

#### 3. 2 Experimental results and discussion

By respectively using three output mirrors w hich have transmittances at  $1331 \text{ nm}$  of  $T_1(20.4\%)$ ,  $T_2$  $( 24. 5\%)$ ,  $T_3( 30. 3\%)$  respectively, we have make the experiment on the  $x$ -axis sample and the  $y$ -axis sample. The results are shown in Fig.  $5a \sim Fig. 5c$ , respectively.  $E_p$  stands for the polarization direction of pumping laser.

From Fig. 5, we know that the output energy of 1331. 0nm fundamental frequency laser is hig hest at  $E_p \parallel z$  and low est at  $E_p \parallel y$ , which is well matched with the polarized transmission spectrum. The highest output energy at 1331. Onm is obtained when the transmissivity is  $T_1$ . For the pump energy of 14mJ, the output energies are 3. 39mJ at  $E_p \parallel z$ , 2. 95mJ at  $E_p \parallel x$  and 1. 85mJ at  $E_p \parallel y$ . The highest slope conversion efficiency is 24. 2%. For the three output mirrors, under three pumping modes  $E_p$  $\parallel x, E_{p} \parallel y$  and  $E_{p} \parallel z$ , the highest and low est pump threshold is 10. 5mJ and 4. 4mJ, respectively . Due to the limited pumping energy, we haven't found the saturated absorption phenomenon in the  $E_{\text{out}}-E_{\text{in}}$ curve. T he measured pulse w idth is 9. 2ns.



Fig. 5 Variation of the fundamental frequency output energy at 1331. 0nm w ith pump energy a-  $E_p$   $\int$  y b-  $E_p$   $\int$  x c-  $E_p$   $\int$  z 1-  $T_1$  = 20. 41% 2-  $T_2$  = 24. 92% 3-  $T_3$  = 30. 33%

T he SFD red laser output has been observed for both  $x$ -axis sample and  $y$ -axis sample. The w aveleng th of red laser measured by WDG5000-1 Grating monochromator is 665. 5nm. Because the  $\gamma$ ax is is more close to the phase matching direction, higher output of SFD red laser can be observed.

### 4 Conclusion

Nd : GdCOB crystal has five quite wide absorption bands w hich are suitable for the xenon lamp pumping. Due to the sensitization of  $\mathrm{Cr}^{3+}$  ,

Cr-Nd-GdCOB crystal is more suitable for the x enon lamp pumping and the pump efficiency is increased by 25%. Moreover, Nd: GdCOB crystal is suitable for dye laser and LD pumping. In addition, Nd-GdCOB crystal has 3 strong fluorescence spectrums, therefore, red, green and blue laser can be obtained by SFD. This ex periment is in prog ress in our laboratory.

### **References**

[ 1] Lu B S, Wang J, Pan H F et al. J A P, 1989,  $3(9)$ :  $413 \sim 416$ -( 下转第 356 页)

, ,



, ,

( 上接第 353 页)

- [2] Mougel F, Aka G, Kahn Harari A et al. J A P, 1998, B67: 533~ 535.
- [ 3] Vivien D, Mougel F, Aka G et al. Laser Phys, 1998, 8( 3) : 759~ 763-
- [4] Mougel F, Aka G, Kahn Harari A et al. Opt Materials,  $1997(8)$ :  $161 - 173$
- [ 5] Zhang H J, Meng X L, Zhu L et al. Opt Commun, 1999, 160: 273  $~1276$
- [6] Aka G, Kahn Harari A, Mougel F. J O S A, 1997, B14(9): 2238  $~2247.$
- [ 7] Wang C Q, Chow Y T, Gambling W A et al. Opt Commun, 2000, 174: 471~ 474-
- [ 8] H ou X Y, Sun Y M, Li Y F et al . Opt & Laser Technol, 2000, 32: 135~ 138-